

# Dressed quantum graphs with optical nonlinearities near the fundamental limits

## Topics

Nonlinearities of quantum wires with finite  $\delta$ -potentials ( $\delta$ -graphs, aka compressed  $\delta$ -atoms and molecules)

Scaling behavior near the fundamental limits (FL)

Origin of the enhanced nonlinearities

## Key results\*

Existence of  $\delta$ -graphs with  $\beta_{\max} \sim 0.71$ ,  $\gamma_{\max} \sim 0.6$ ,  $\gamma_{\min} \sim -0.15$

Three-levels required for  $\beta$  near its maximum, & for  $\gamma$  near its minimum

Four levels required for  $\gamma$  near its maximum

\* *all hyperpolarizabilities are x-diagonal and normalized to the FL*

**Rick Lytel**

August 7, 2013

# Reading list

“Analytical solution of the compressed, one-dimensional delta atom via quadratures and exact, absolutely convergent periodic-orbit expansions,” R. Blümel, *Journal of Physics A: Mathematical and General* 39, 8257 (2006).

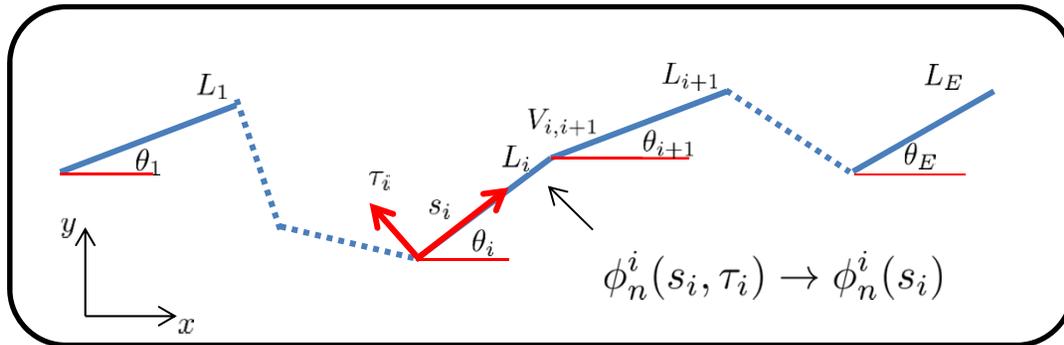
“Scaling and universality in nonlinear optical quantum graphs containing star motifs,” R. Lytel, S. Shafei, & M. Kuzyk, *arXiv:1305.4334* (2013).

“Influence of geometry and topology of quantum graphs on their nonlinear optical properties,” Rick Lytel, Shoresh Shafei, Julian H. Smith, and Mark G. Kuzyk, *Phys. Rev. A* 87, 043824 (2013).

“Sum Rules and Scaling in Nonlinear Optics,” M.G. Kuzyk, J. Perez-Moreno, and S. Shafei, *Phys. Rep.* 529, 297-398 (2013).

“Geometry-controlled nonlinear optical response of quantum graphs,” S. Shafei, R. Lytel, and M.G. Kuzyk, *J. Opt. Sci. Am. B* 29, 3419 (2012).

# Quantum mechanics of **bare** graphs



Extreme confinement along  $\tau$

Electron dynamics along  $s$

Self-adjoint Hamiltonian

**Zero potential everywhere**

Eigenfunctions are unions of edge functions having same eigenvalues

$$H\psi_n(s) = E_n\psi_n(s) \quad \psi_n(s) = \cup_{i=1}^E \phi_n^i(s_i) \quad H\phi_n^i = E_n\phi_n^i$$

Edge functions: continuity and conservation of probability flux

$$\phi_n^i = A_n^{(i)} \sin k_n s_i + B_n^{(i)} \cos k_n s_i \quad \phi_n^i = \phi_n^{i+1} \text{ at } V_{i,i+1} \quad \sum_{i=1}^d \phi_n^i = 0 \text{ at } V_{i,i+1}$$

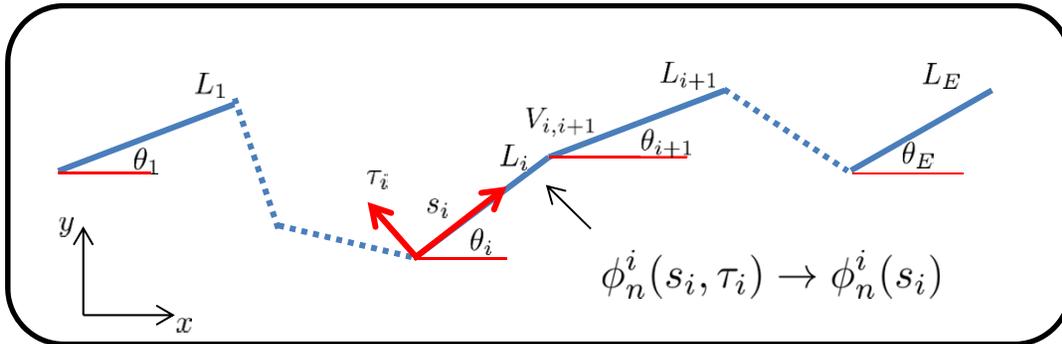
**'degree' of vertex**

$d$

Homogeneous, coupled linear equations for amplitudes

$$\sum_{i=1}^{2E} a_n^{ij} C_n^{(i)} = 0, \quad j = 1, 2, 3 \dots 2E, \quad C_n^{(1)} = A_n^{(1)}, \quad C_n^{(2)} = B_n^{(1)} \dots$$

# Quantum mechanics of $\delta$ -dressed graphs



Extreme confinement along  $\tau$

Electron dynamics along  $s$

Self-adjoint Hamiltonian

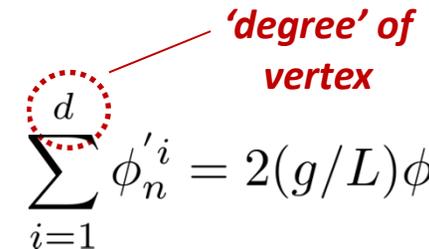
$V = (g/L)\delta(s)$  at a vertex

Eigenfunctions are unions of edge functions having same eigenvalues

$$H\psi_n(s) = E_n\psi_n(s) \quad \psi_n(s) = \cup_{i=1}^E \phi_n^i(s_i) \quad H\phi_n^i = E_n\phi_n^i$$

Edge functions: continuity and conservation of probability flux

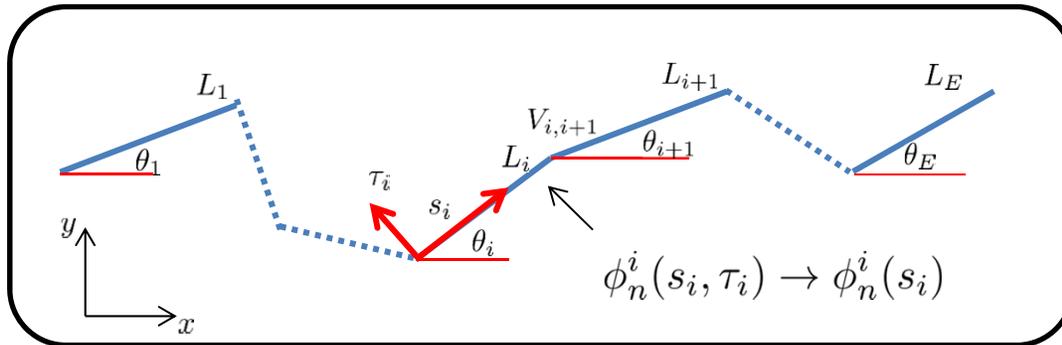
$$\phi_n^i = A_n^{(i)} \sin k_n s_i + B_n^{(i)} \cos k_n s_i \quad \phi_n^i = \phi_n^{i+1} \text{ at } V_{i,i+1} \quad \sum_{i=1}^d \phi_n^i = 2(g/L)\phi_n(a)$$



Homogeneous, coupled linear equations for amplitudes

$$\sum_{i=1}^{2E} a_n^{ij} C_n^{(i)} = 0, \quad j = 1, 2, 3 \dots 2E, \quad C_n^{(1)} = A_n^{(1)}, \quad C_n^{(2)} = B_n^{(1)} \dots$$

# Quantum mechanics of graphs (cont'd)



Extreme confinement along  $\tau$

Electron dynamics along  $s$

Self-adjoint Hamiltonian

Bare or dressed edges

Determinant of coefficients vanishes yielding secular equation for eigenvalues

$$\det a_n^{ij} = 0 \quad f_{sec}(k_n; L_i, \theta_i) = 0 \Rightarrow E_n = k_n^2/2$$

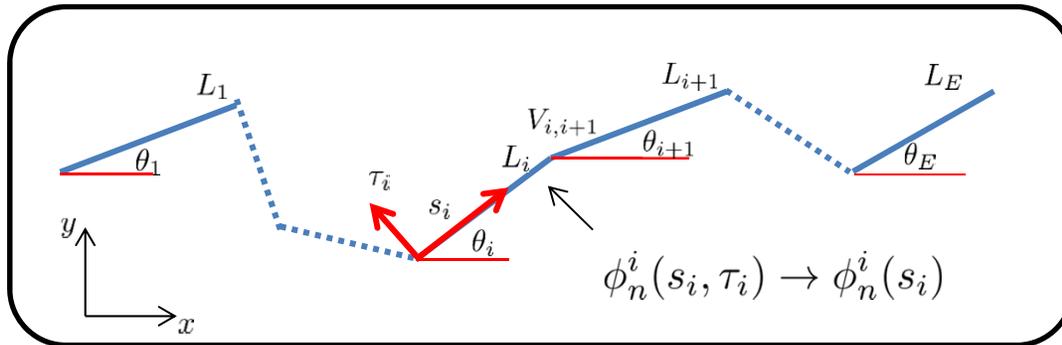
Continuity equations for amplitudes, plus normalization yields amplitudes

$$\phi_n^i = A_n^{(i)} \sin k_n s_i + B_n^{(i)} \cos k_n s_i$$

Transition moments are now calculable

$$r_{nm} = \sum_{i=1}^E \int ds_i r(s_i) \phi_n^{*i}(s_i) \phi_m^i(s_i)$$

## Aside: boundary condition equation counting



$E$  edges

$V_E$  external vertices

$V_I$  internal vertices

$V = V_E + V_I$  vertices

Union of edge functions has  $2E$  unknown amplitudes

$$\phi_n^i = A_n^{(i)} \sin k_n s_i + B_n^{(i)} \cos k_n s_i, \quad i = 1, 2, \dots, E$$

1. There are exactly  $V_E$  amplitude boundary condition eqns on terminated edges
2. There are exactly  $V_I$  flux boundary condition eqns on internal vertices
3. This provides exactly  $V$  boundary condition eqns (so far)

Now, let  $d^k$  = degree of each internal vertex, where degree = #edges connected to it. There are exactly  $d^k - 1$  amplitude boundary condition eqns at vertex  $k$ . Sum over internal vertices. The sum over degrees equals the total # of edges that are internally connected, and is equal to  $2E - V_E$ . The sum of 1 over internal edges is  $V_I$ . So we get:

4. There are exactly  $2E - V_E - V_I = 2E - V$  amplitude boundary eqns at internal vertices. QED

# First hyperpolarizability tensor

$$\beta_{ijk} = -\frac{e^3}{2} \sum'_{n,m} \frac{1}{E_{n0}E_{m0}} \left( r_{0n}^i \bar{r}_{nm}^j r_{m0}^k + r_{0n}^i \bar{r}_{nm}^k r_{m0}^j + r_{0n}^k \bar{r}_{nm}^i r_{m0}^j + h.c. \right)$$

**Fully symmetric 3<sup>rd</sup> rank Cartesian tensor**

$$\begin{aligned} \beta_{xxx}(\theta) &= \beta_{xxx} \cos^3 \theta + 3\beta_{xxy} \cos^2 \theta \sin \theta \\ &+ 3\beta_{xyy} \cos \theta \sin^2 \theta + \beta_{yyy} \sin^3 \theta \end{aligned}$$

**Four independent components relating lab and graph frames**

$$|\beta| = \sqrt{\beta_{xxx}^2 + 3\beta_{xxy}^2 + 3\beta_{xyy}^2 + \beta_{yyy}^2}$$

**Tensor norm (invariant under rotations of graph)**

$$\beta_{max} = 3^{1/4} \left( \frac{e\hbar}{\sqrt{m}} \right)^3 \left( \frac{N^{3/2}}{E_{10}^{7/2}} \right) \longrightarrow -1 \leq \beta_{int} \equiv \frac{\beta}{\beta_{max}} \leq 1$$

**Normalization for scale-invariant calculations**

# Second hyperpolarizability tensor

$$\gamma_{ijkl} = (1/6)P_{ijkl} \left( \sum_{n,m,l} \frac{r_{0n}^i \bar{r}_{nm}^j \bar{r}_{ml}^k r_{l0}^l}{E_{n0} E_{m0} E_{l0}} - \sum_{n,m} \frac{r_{0n}^i r_{n0}^j r_{0m}^k r_{m0}^l}{E_{n0}^2 E_{m0}} \right)$$

**Fully symmetric 4<sup>th</sup> rank Cartesian tensor**

$$\begin{aligned} \gamma_{xxxx}(\theta) &= \gamma_{xxxx} \cos^4 \theta + 4\gamma_{xxxy} \cos^3 \theta \sin \theta \\ &+ 6\gamma_{xxyy} \cos^2 \theta \sin^2 \theta + 4\gamma_{xyyy} \cos \theta \sin^3 \theta \\ &+ \gamma_{yyyy} \sin^4 \theta \end{aligned}$$

**Five independent components relating lab and graph frames**

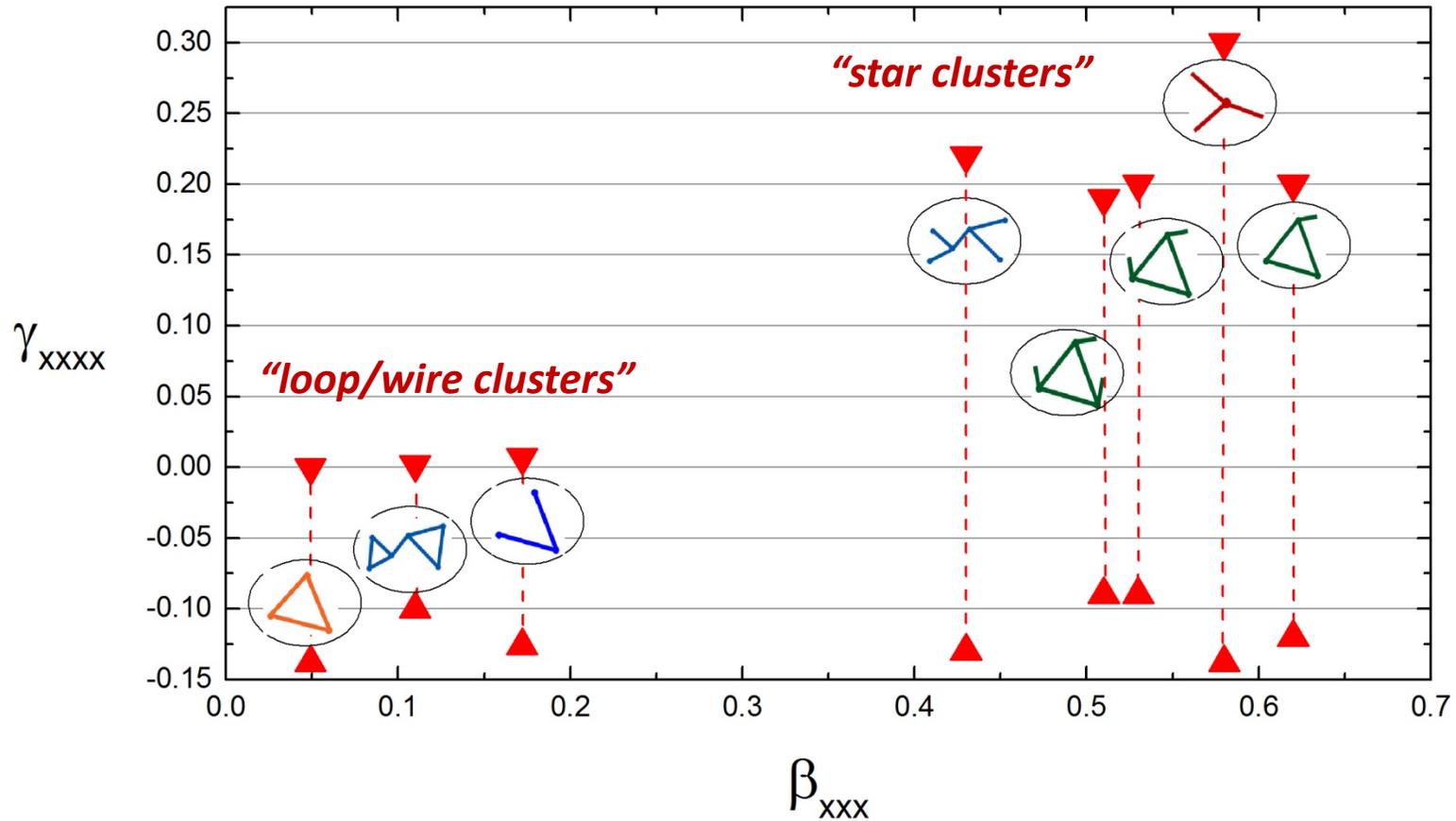
$$|\gamma| = \sqrt{\gamma_{xxxx}^2 + 4\gamma_{xxxy}^2 + 6\gamma_{xxyy}^2 + 4\gamma_{xyyy}^2 + \gamma_{yyyy}^2}$$

**Tensor norm (invariant under rotations of graph)**

$$\gamma_{max} = 4 \left( \frac{e^4 \hbar^4}{m^2} \right) \left( \frac{N^2}{E_{10}^5} \right) \longrightarrow -1/4 \leq \gamma_{int} \equiv \frac{\gamma}{\gamma_{max}} \leq 1$$

**Normalization for scale-invariant calculations**

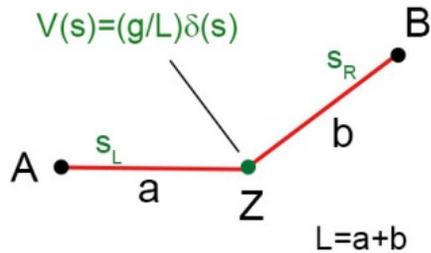
# Topological clustering and enhancement of $\beta$ and $\gamma$



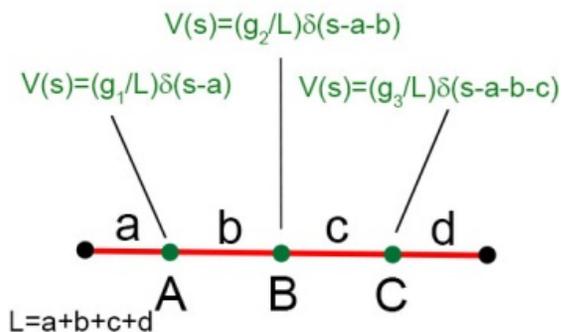
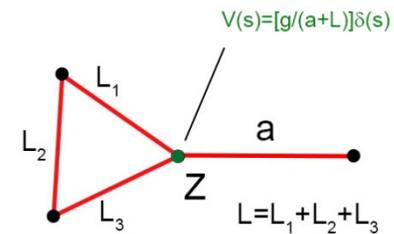
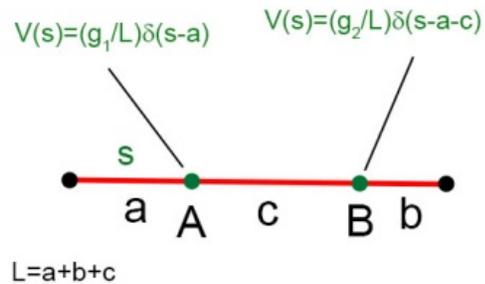
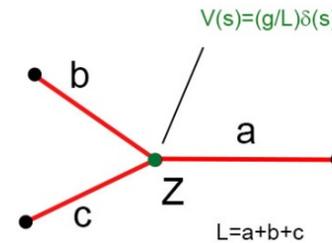
Hyperpolarizability ranges still fall short of (potential optimization) limits

# $\delta$ -dressed quantum graphs

*“dressed wires”*

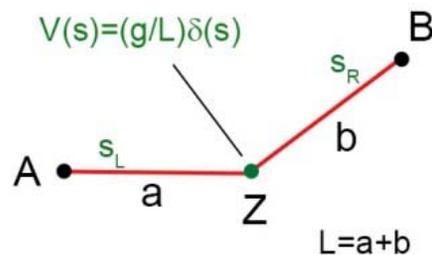


*“dressed stars”*



Presence of finite  $\delta$  potential re-arranges mode shapes  
Might we create the optimum states for  $\max \beta, \gamma$ ?

# Compressed delta atom motif



$$\phi_L(s) = \frac{Z \sin k(a - s_L) + A \sin k s_L}{\sin k a}$$

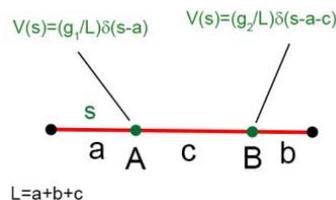
$$\phi_R(s) = \frac{Z \sin k(b - s_R) + B \sin k s_R}{\sin k b}$$

**Secular function defined by flow of probability on the graph:**

$$Z F_\delta = A \sin k b + B \sin k a \quad \omega = 2a/L - 1, -1 \leq \omega \leq 1$$

$$F_\delta \equiv F_\delta(g; \omega, kL) = -(1/kL) [g(\cos kL - \cos \omega kL) - kL \sin kL]$$

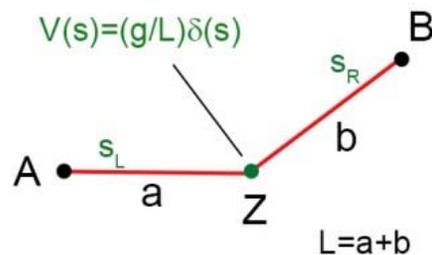
**Overlapping motifs, with matched endpoints, yields secular functions for larger graphs**



$$A F_\delta(g_1; a, c) = B \sin k a$$

$$B F_\delta(g_2; b, c) = A \sin k b$$

## Compressed delta atom (A=B=0)



$$\phi_L(s) = \frac{Z \sin k(a - s_L)}{\sin ka}$$

$$\phi_R(s) = \frac{Z \sin k(b - s_R)}{\sin kb}$$

### Secular equation determines eigenvalues for graph

$$0 = F_\delta(g; \omega, kL) = -(1/kL) [g(\cos kL - \cos \omega kL) - kL \sin kL], \quad E \geq 0$$

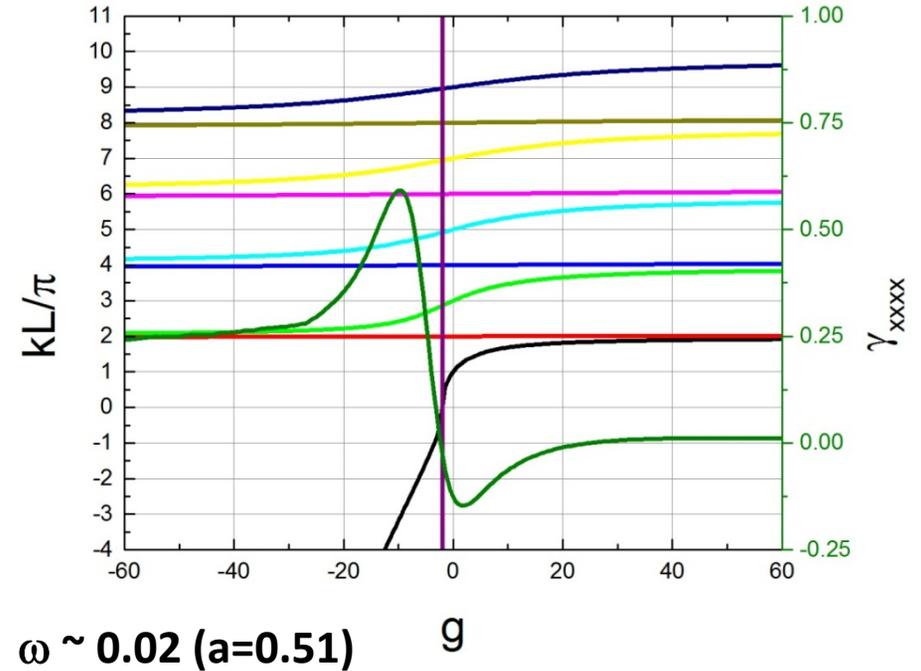
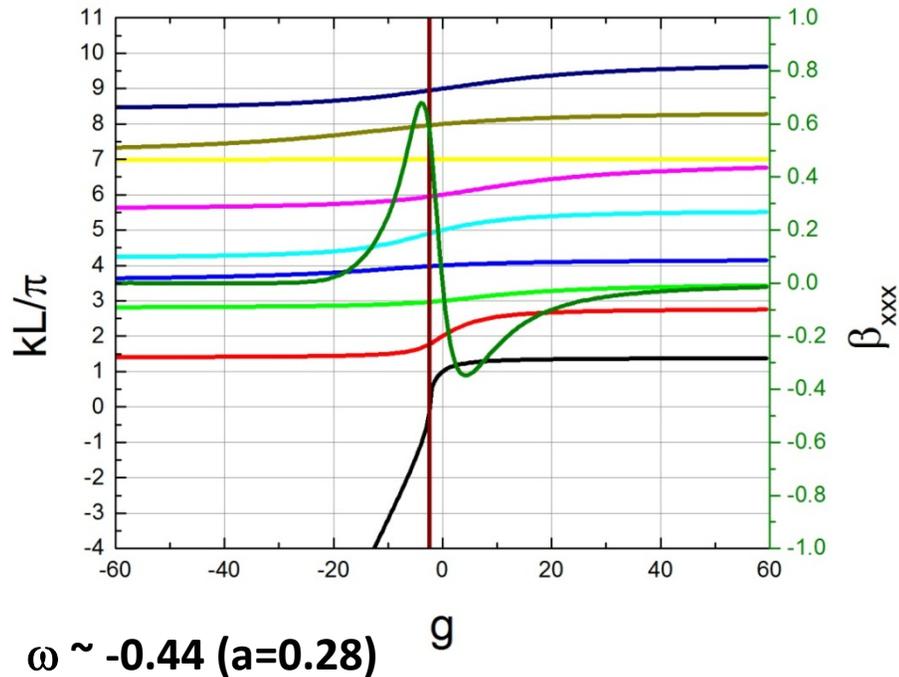
$$0 = F_\delta(g; \omega, \kappa L) = (i/\kappa L) [g(\cosh \kappa L - \cosh \omega \kappa L) + \kappa L \sinh \kappa L], \quad E < 0$$

$$E_n = k_n^2/2, \quad \hbar = m = 1, \quad k = i\kappa, \quad E < 0$$

### Finite potential allows a single negative energy state for negative \$g < g\_c\$

$$-ixF_\delta(g; \omega, \kappa L) = x^2 \left[ 1 - \frac{g}{g_c} \right] + O(x^4), \quad x \equiv \kappa L, \quad g_c = 2/(\omega^2 - 1)$$

# Spectra and hyperpolarizabilities for finite vertex potentials

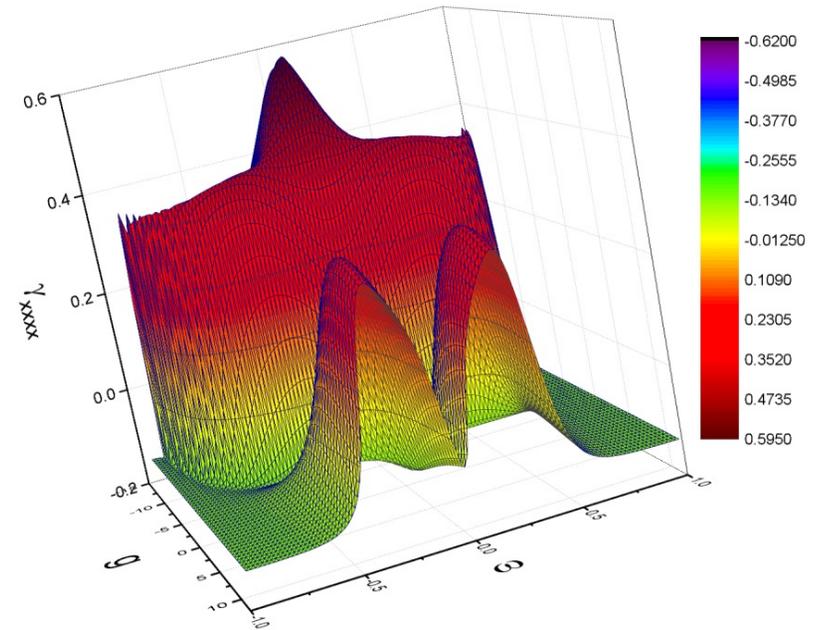
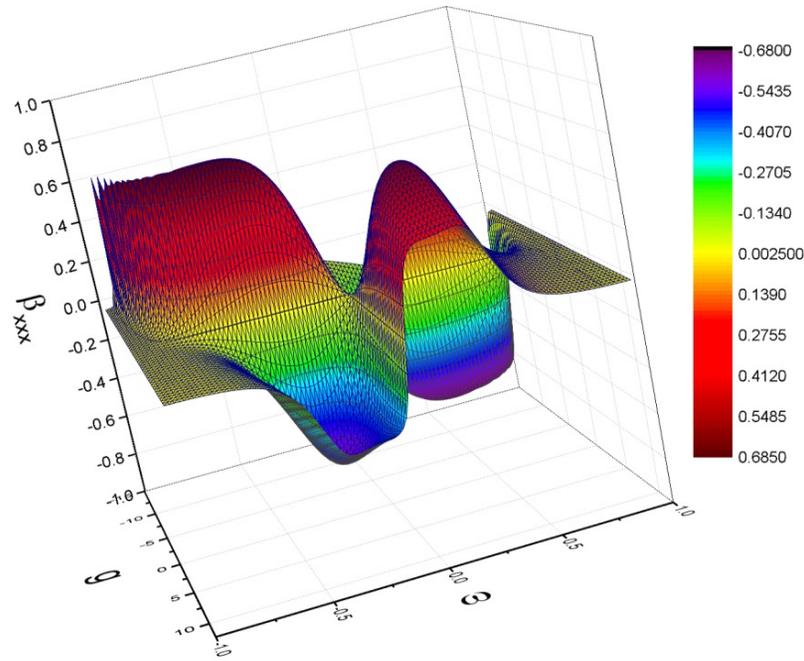


Positive eigenvalues lie between fixed root boundaries at  $n\pi/L$

Topologically-induced spectrum shift when negative energy state appears

Shift induces giant enhancement of nonlinearities. **WHY?**

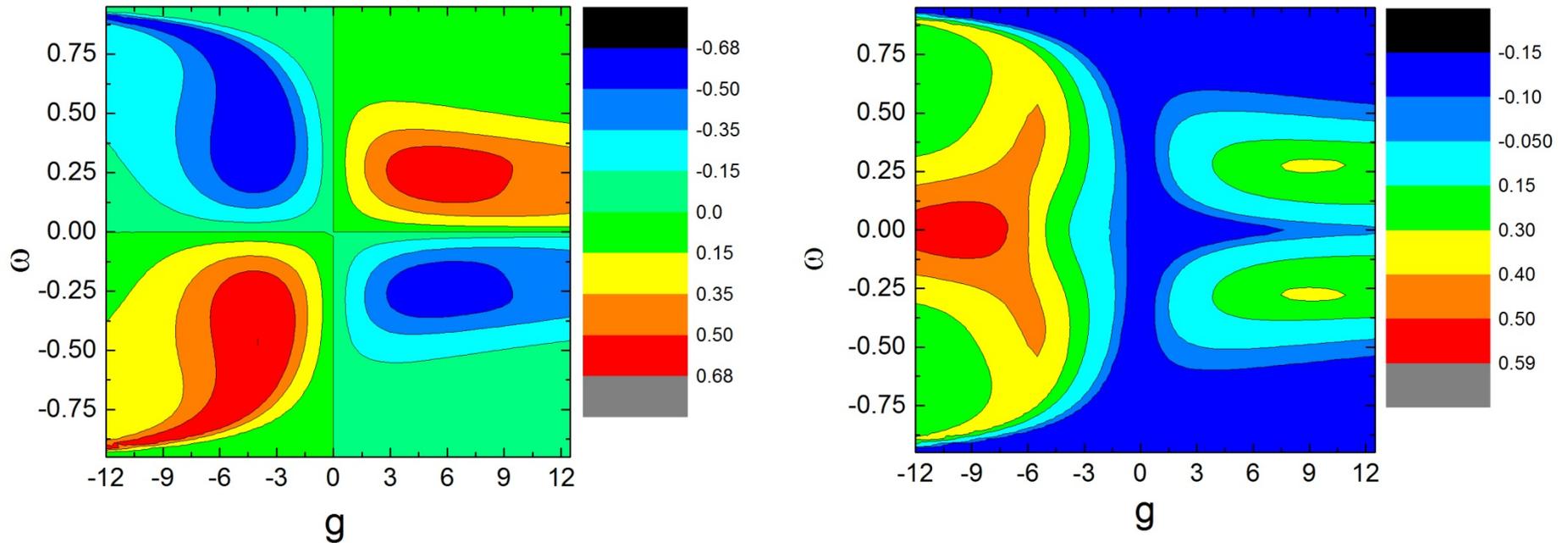
# 3D views of ranges of hyperpolarizabilities



$\gamma$  saturates at -0.126 when the  $\delta$  function is moved to either edge (bare wire)

$\beta$  vanishes when the  $\delta$  function is moved to either edge or to the center

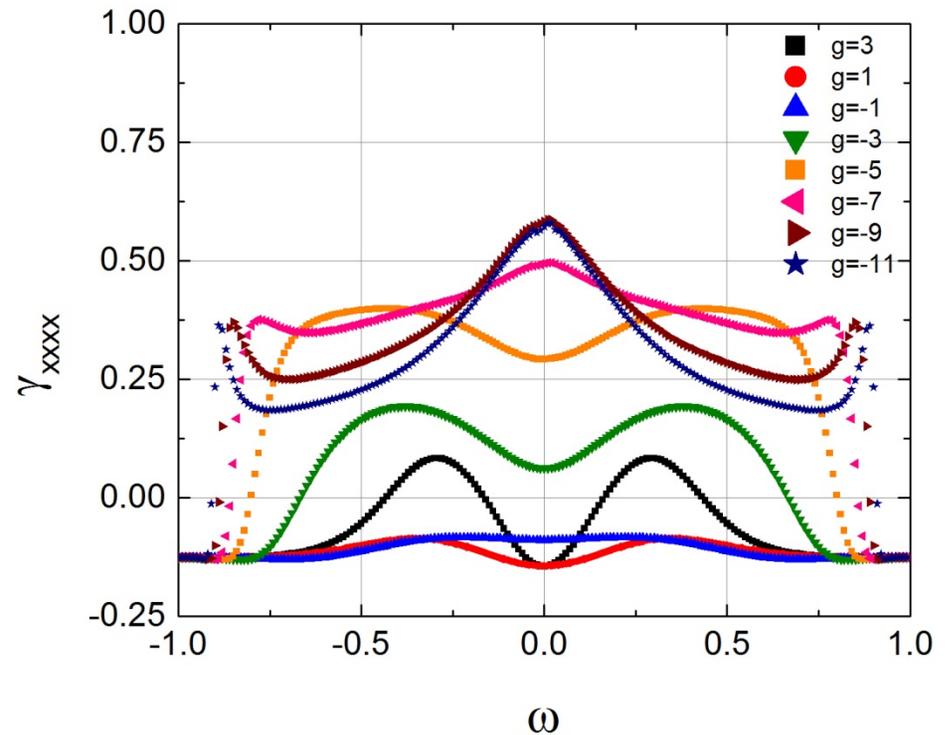
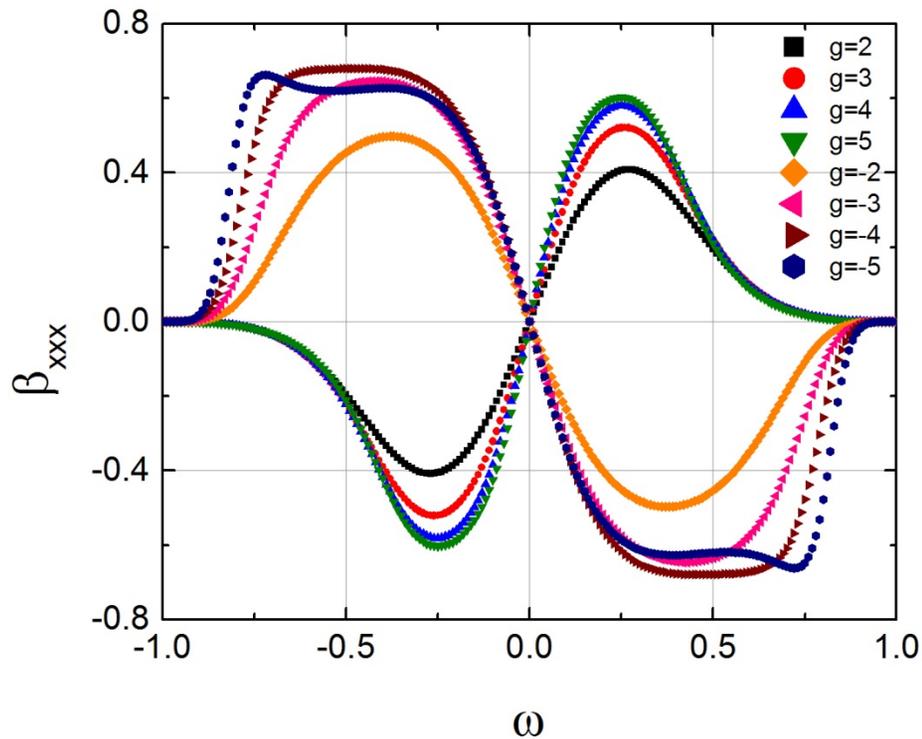
# Contour views of ranges of hyperpolarizabilities



$\gamma$  saturates at -0.126 when the  $\delta$  function is moved to either edge (bare wire)

$\beta$  vanishes when the  $\delta$  function is moved to either edge or to the center

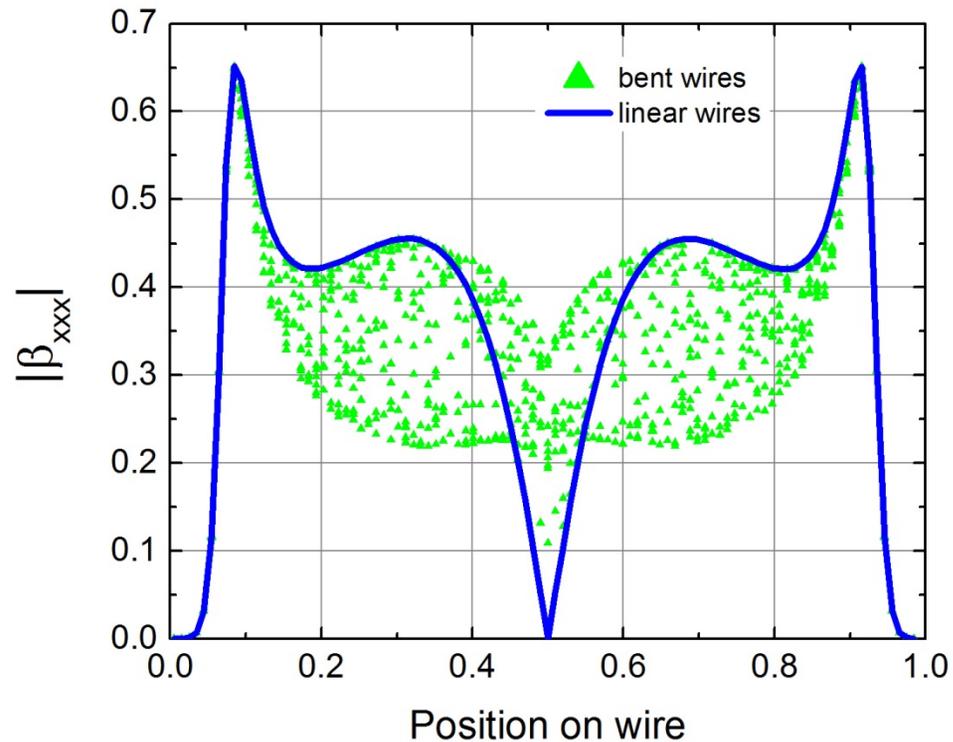
# Ranges of hyperpolarizabilities across the atom



$\gamma$  saturates at -0.126 when the  $\delta$  function is moved to either edge (bare wire)

$\beta$  vanishes when the  $\delta$  function is moved to either edge or to the center

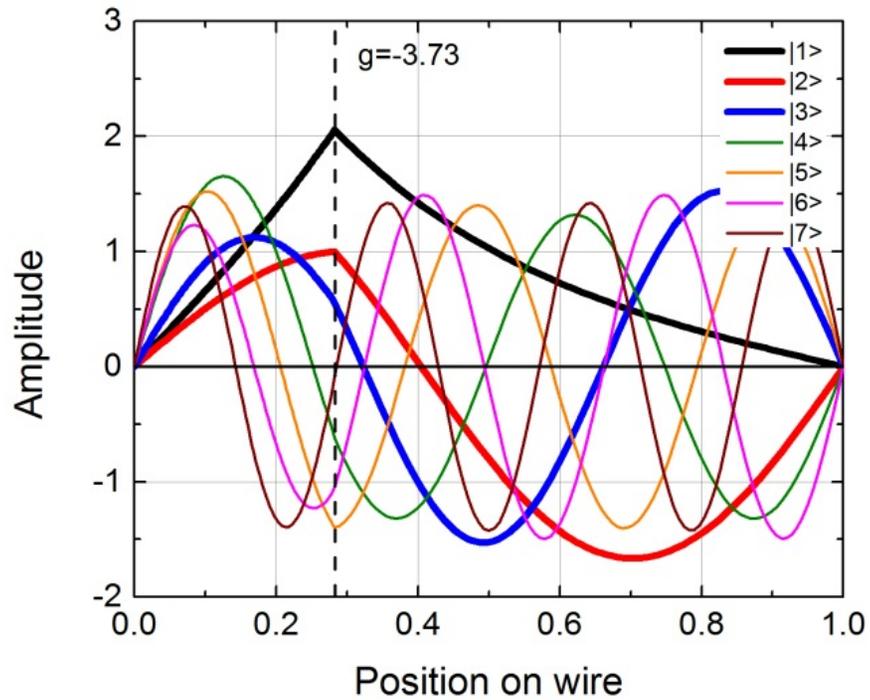
# Effect of bending one edge relative to the other



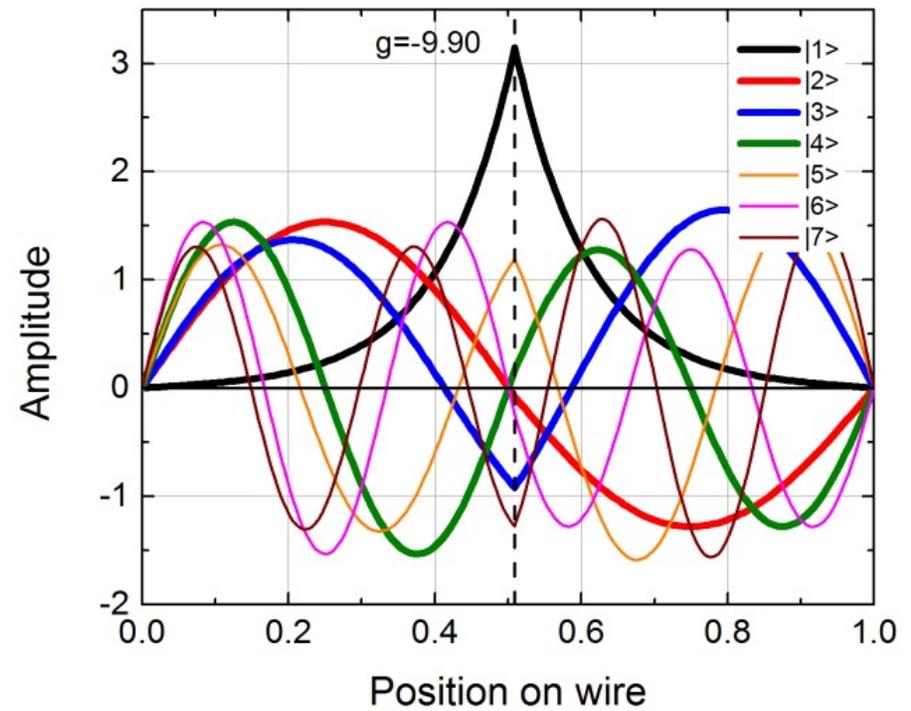
$g = -7$

Linear graph has the largest  $\beta$  except near the center of the graph

# First seven eigenstates for best $\beta$ and $\gamma$



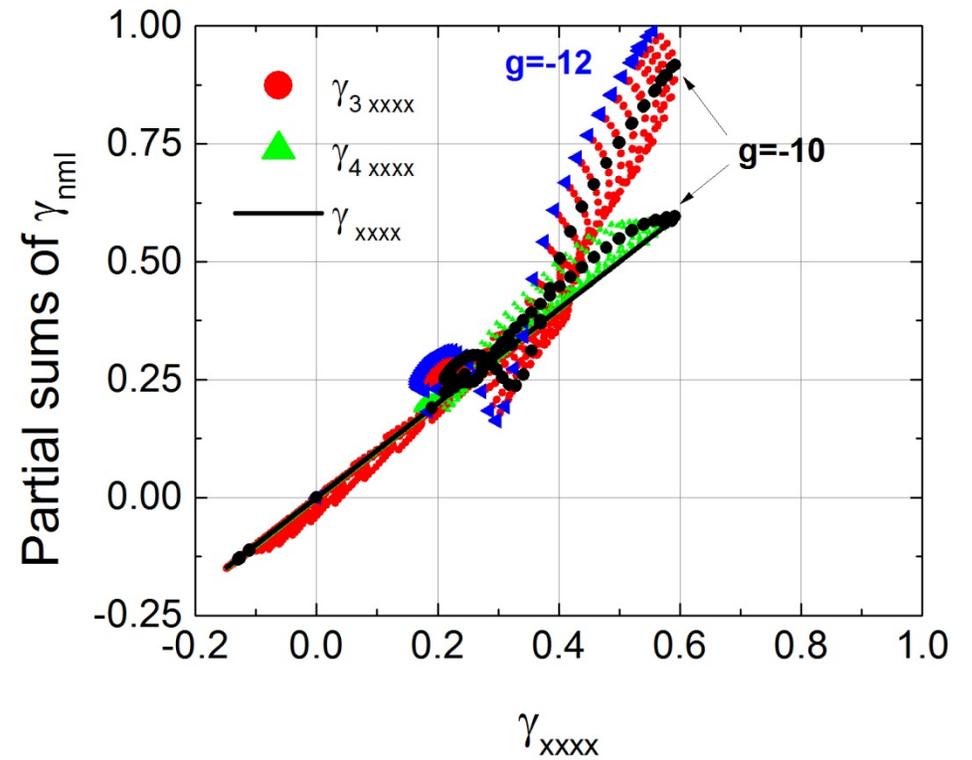
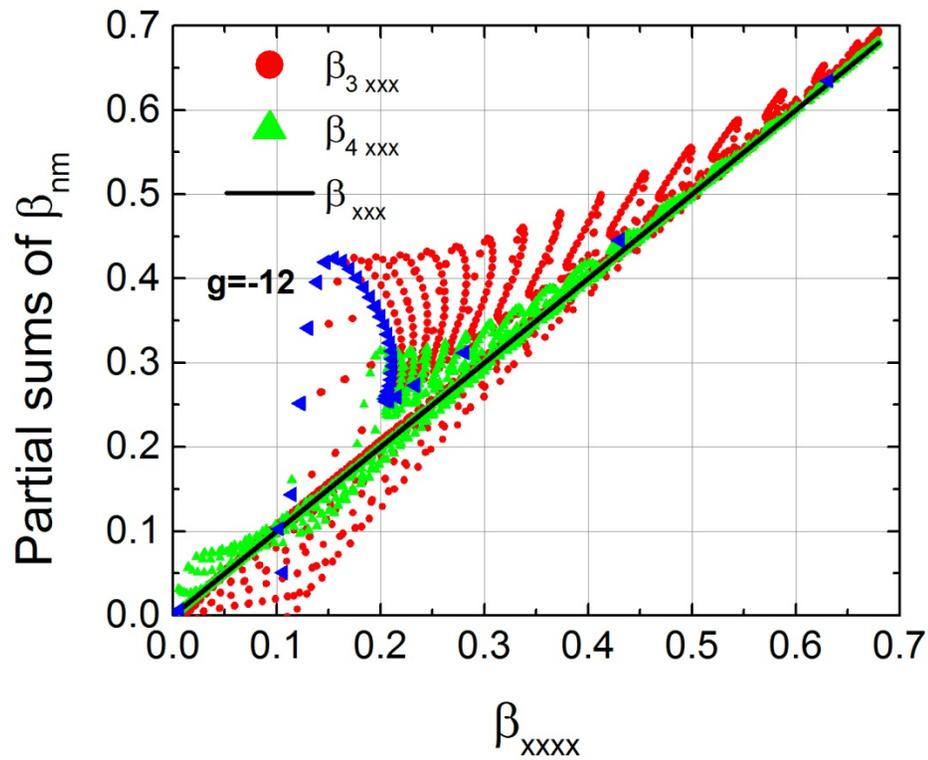
$\beta \sim 0.68$



$\gamma \sim 0.58$

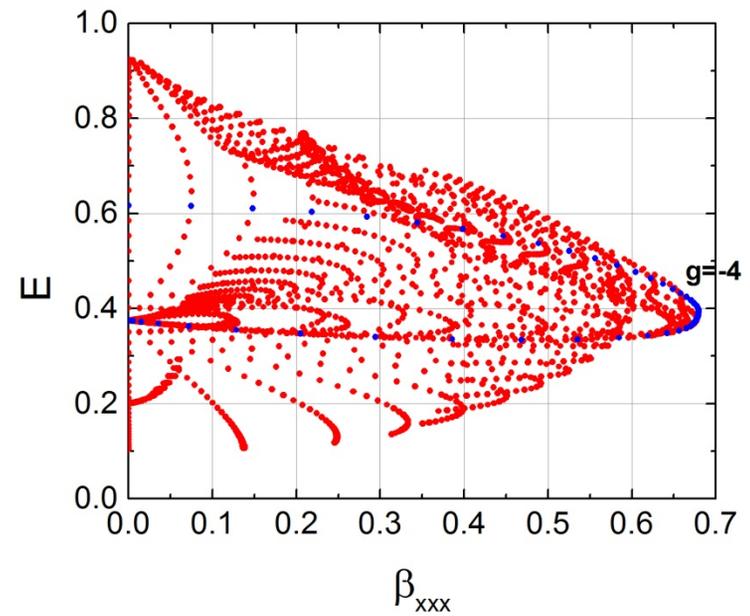
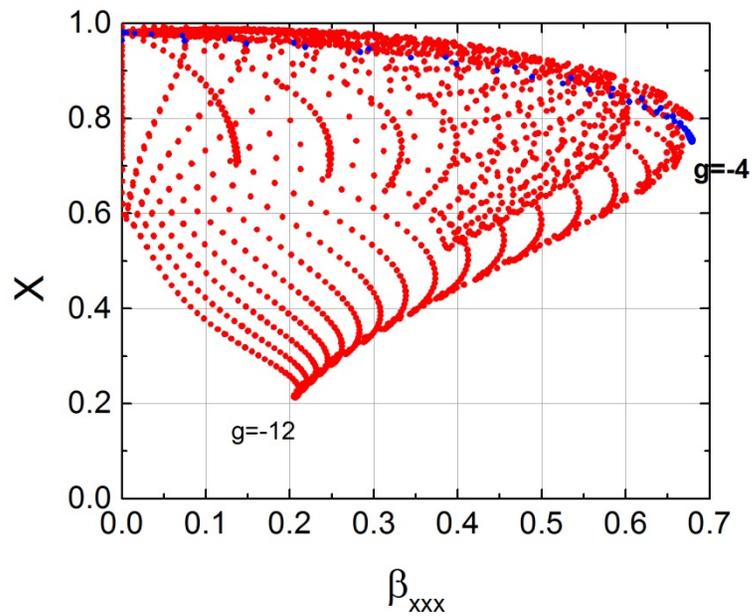
Exquisite tuning of eigenstate overlap for transition moments

# Three- and four-level Ansatz



$g=-12$  to  $+12.5$  in steps of 0.5, each with 200 positions along the wire

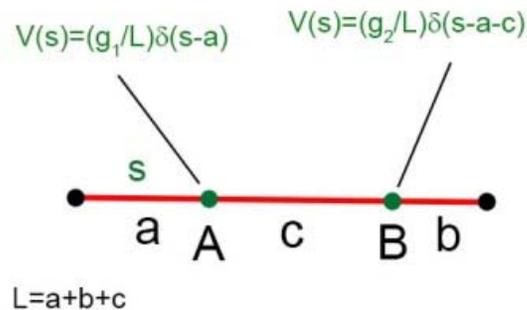
# Convergence to universal scaling values for X and E



Normalized dipole moment  $X$  asymptotes to 0.79 at maximum  $\beta$

Normalized energy ratio  $E$  asymptotes to 0.39 at maximum  $\beta$

# Compressed 2-delta molecule



$$\begin{aligned}\phi_a(s_a) &= \frac{A \sin ks_a}{\sin ka} \\ \phi_c(s_c) &= \frac{A \sin k(c - s_c) + B \sin ks_s}{\sin kc} \\ \phi_b(s_b) &= \frac{B \sin k(s_b - b)}{\sin kb}\end{aligned}$$

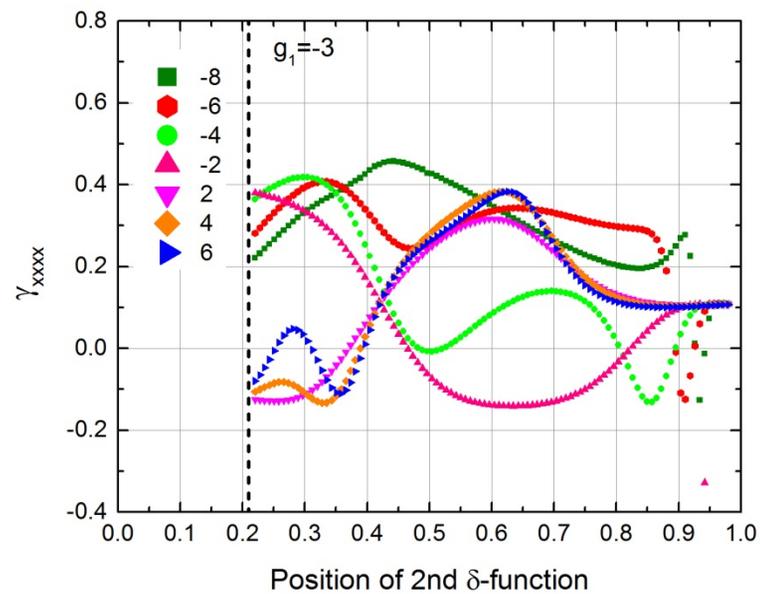
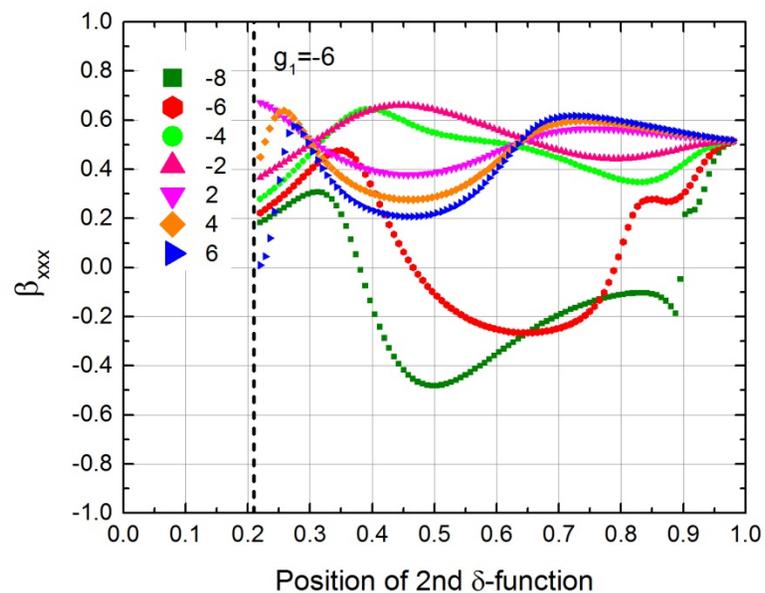
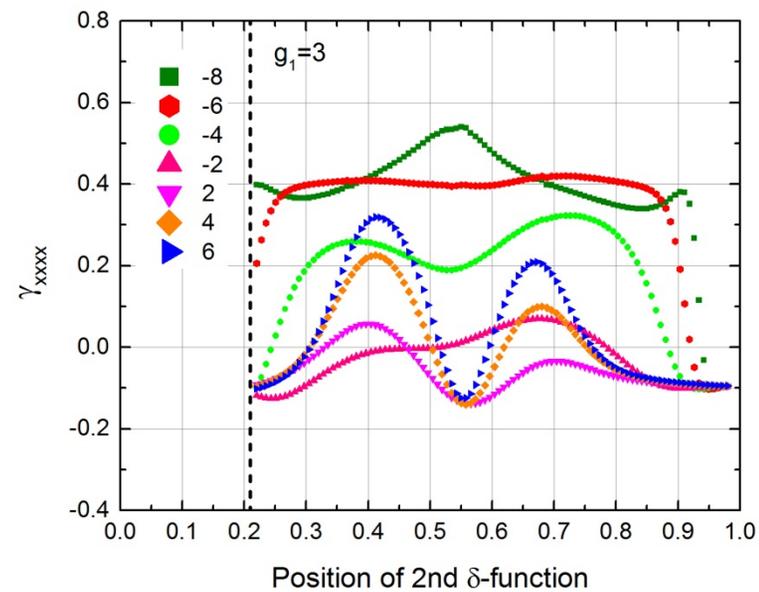
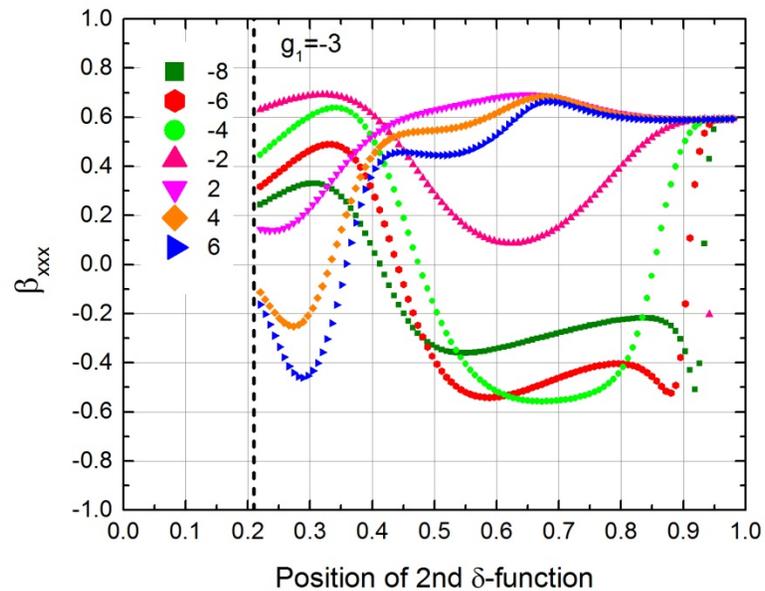
**Overlapping motifs, with matched endpoints, yields secular functions for larger graphs**

$$\begin{aligned}ZF_\delta &= A \sin kb + B \sin ka \\ AF_\delta(g_1; a, c) &= B \sin ka \\ BF_\delta(g_2; b, c) &= A \sin kb\end{aligned}$$

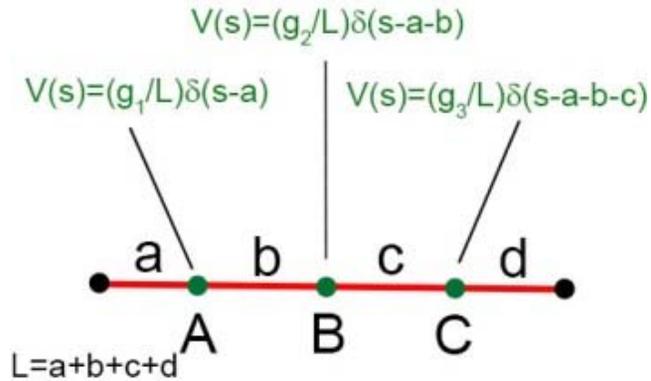
**Secular equation determines eigenvalues for graph (two negative energy states possible)**

$$\begin{aligned}F_{2\delta}(g_1, g_2; a, b, c) &= F_\delta(g_1; a, c)F_\delta(g_2; b, c) - \sin ka \sin kb = 0 \\ F_\delta(g_1; \omega_1, kL) &= -(1/kL) [g_1(\cos kL_1 - \cos \omega_1 kL_1) - kL \sin kL_1] \\ F_\delta(g_2; \omega_2, kL) &= -(1/kL) [g_2(\cos kL_2 - \cos \omega_2 kL_2) - kL \sin kL_2]\end{aligned}$$

# Hyperpolarizabilities across wire



# Compressed 3-delta molecule



## Overlapping motifs

$$ZF_{\delta} = A \sin kb + B \sin ka$$

$$AF_{\delta}(g_1; a, b) = B \sin ka$$

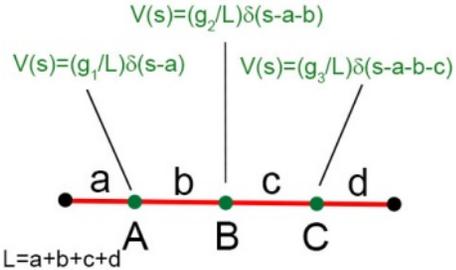
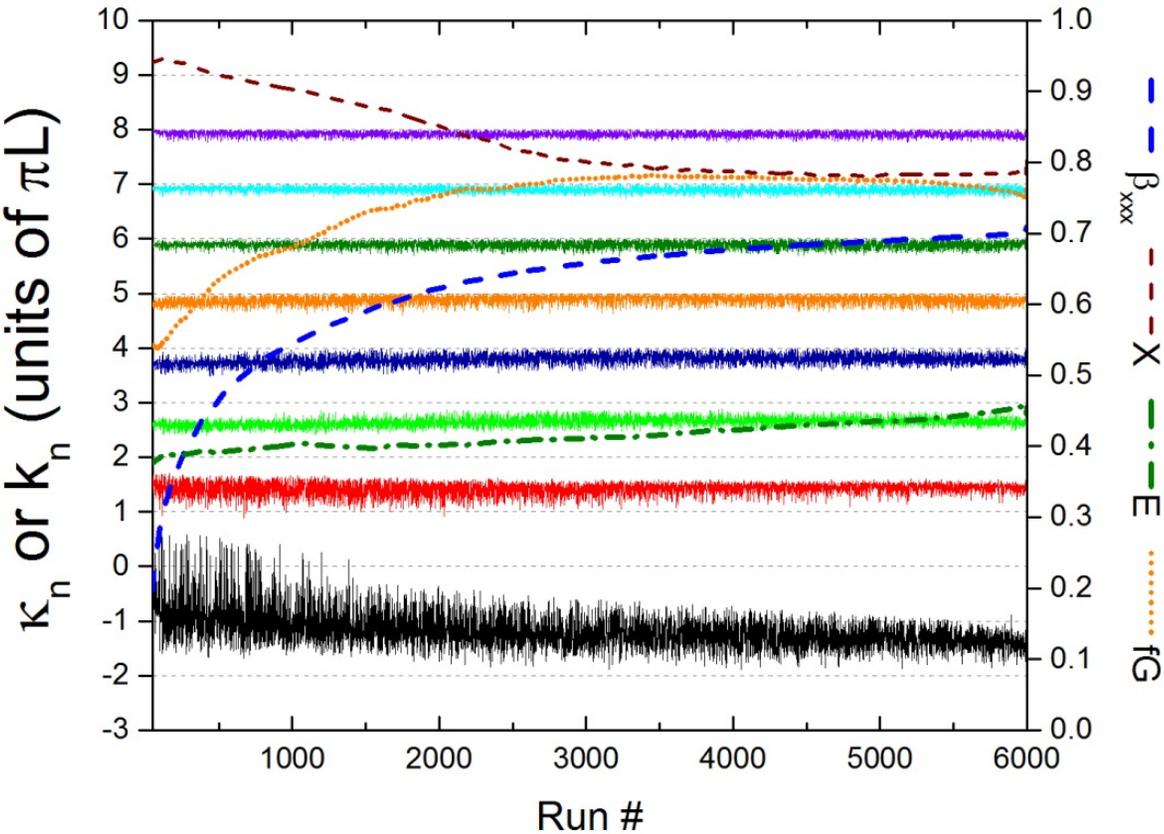
$$BF_{\delta}(g_2; b, c) = A \sin kc + C \sin kb$$

$$CF_{\delta}(g_3; c, d) = B \sin kd$$

Secular equation determines eigenvalues for graph (three negative energy states possible)

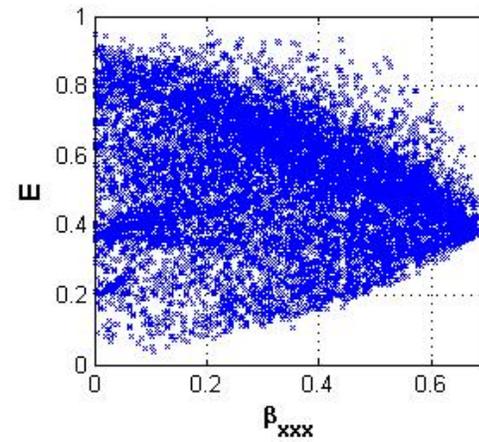
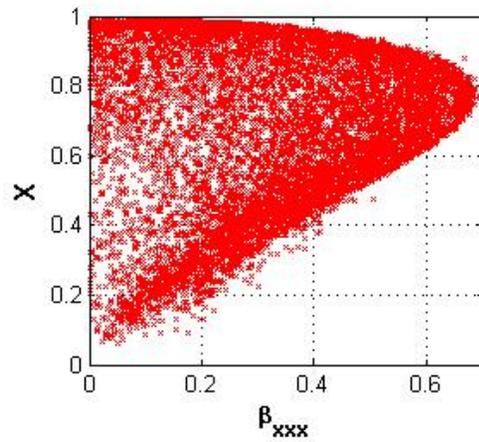
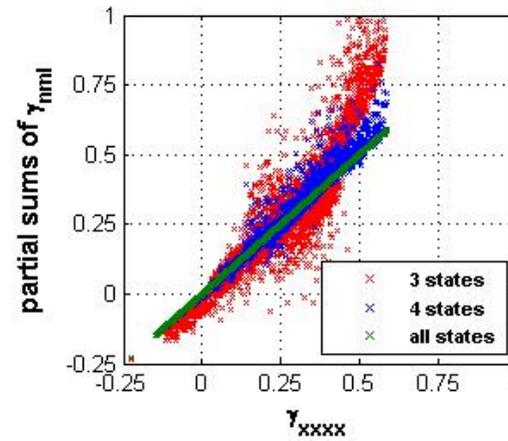
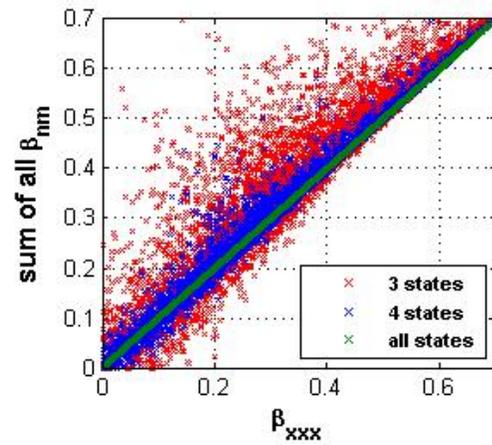
$$\begin{aligned} F_{3\delta}(g_1, g_2, g_3; a, b, c, d) &= F_{\delta}(g_1; a, b)F_{\delta}(g_2; b, c)F_{\delta}(g_3; c, d) \\ &- F_{\delta}(g_1; a, b) \sin kb \sin kd - F_{\delta}(g_3; c, d) \sin ka \sin kc \\ &= 0 \end{aligned}$$

# Spectral variation for random $\delta$ strengths/locations

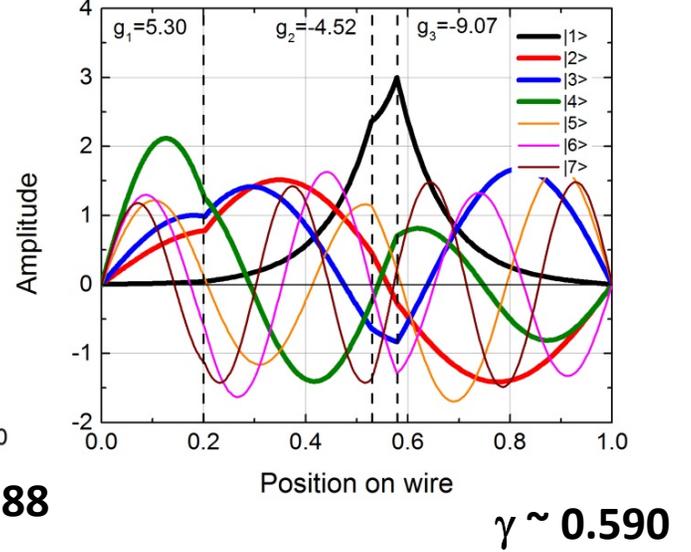
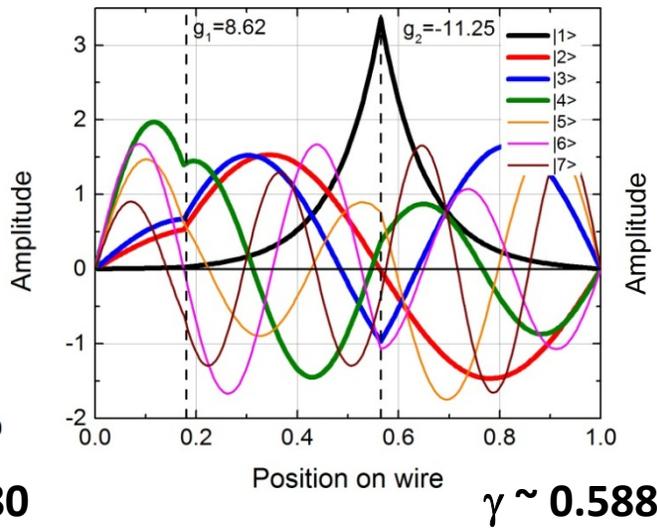
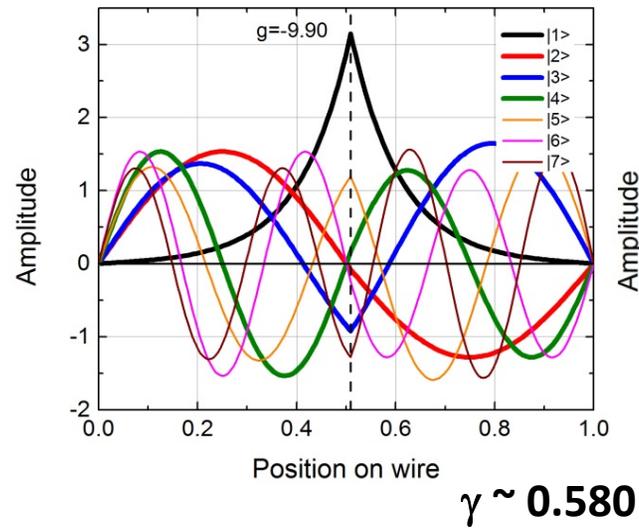
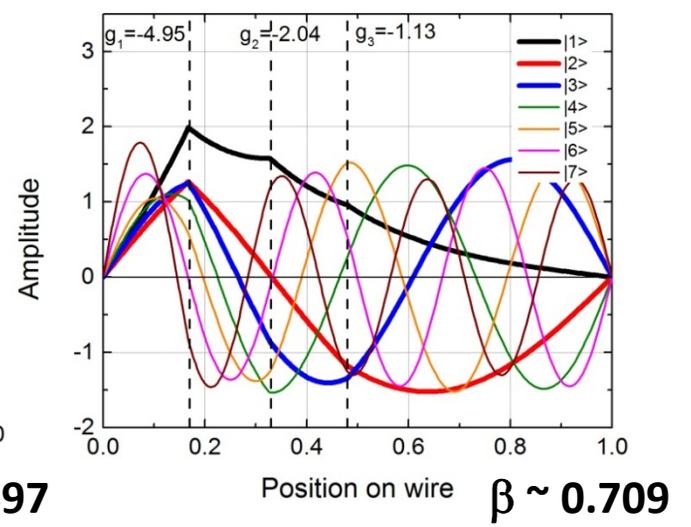
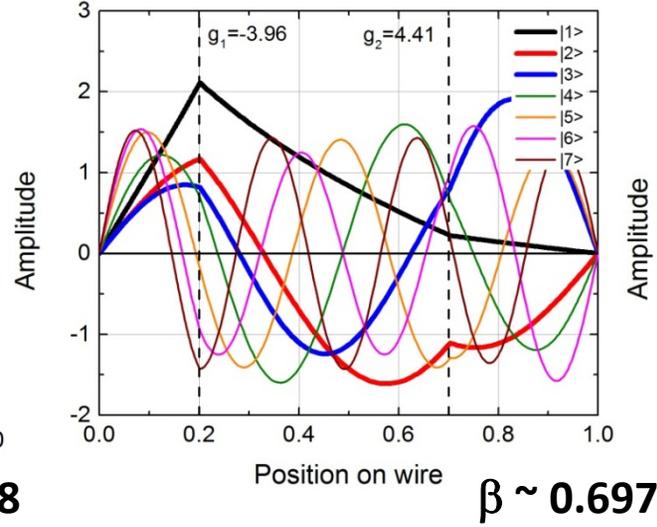
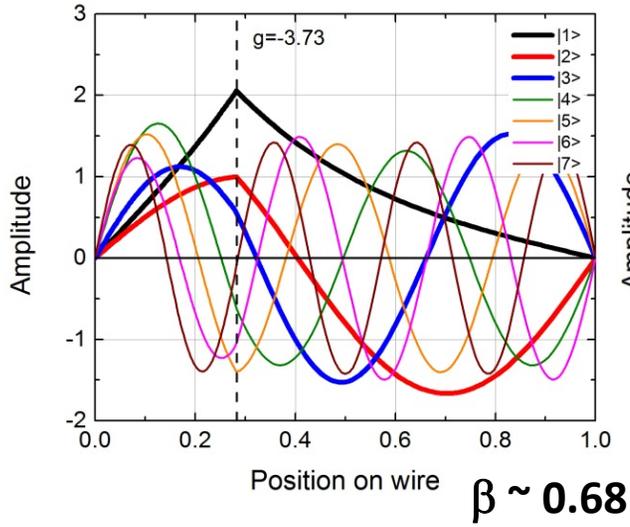


# 10,000 3-delta, random positions

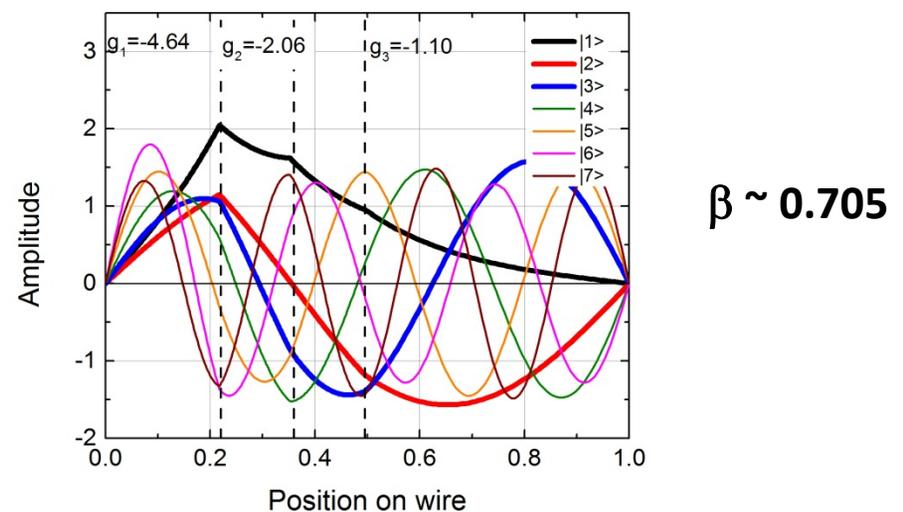
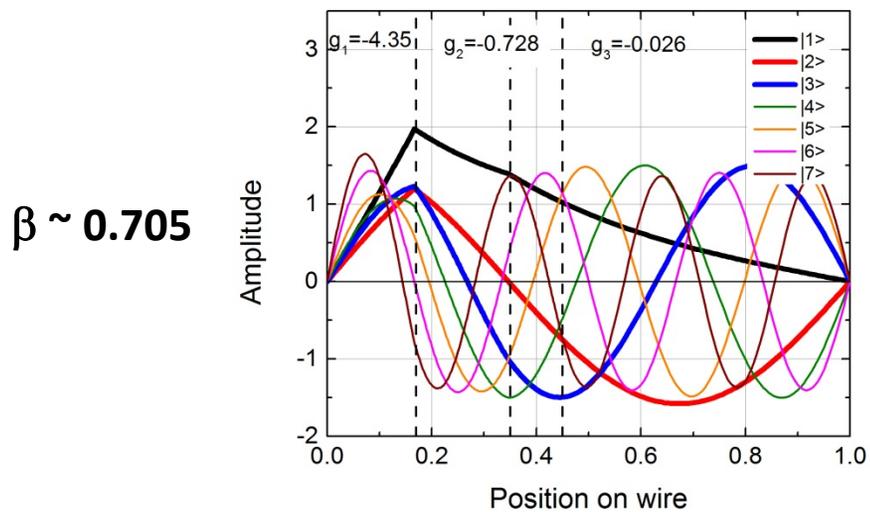
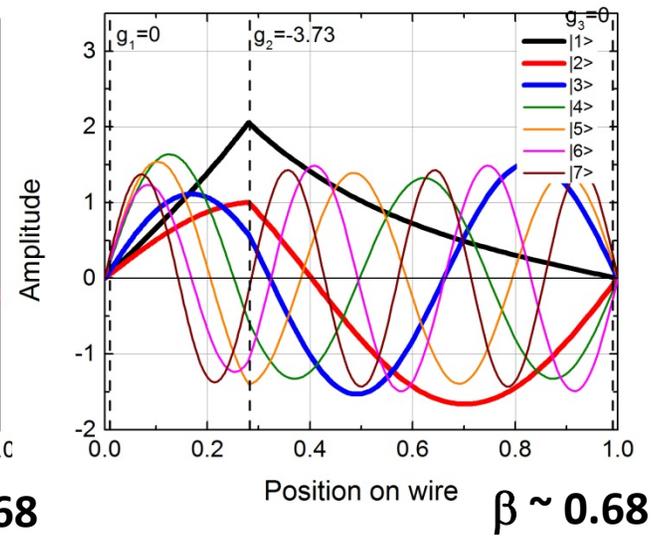
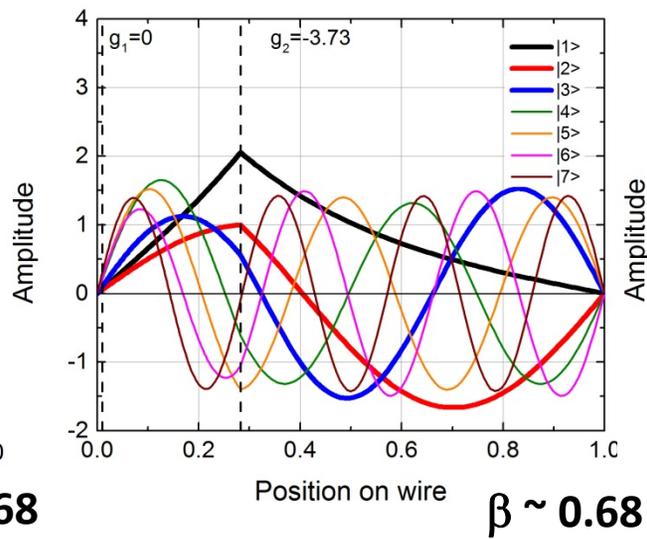
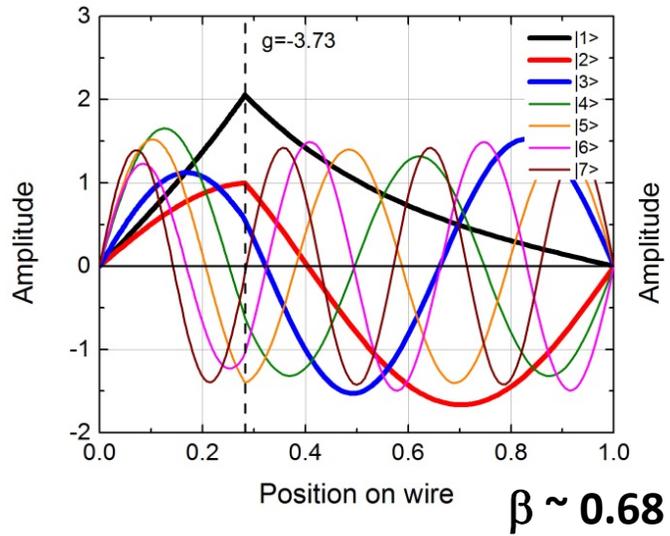
$$|g_1| < 8, |g_2| < 8, |g_3| < 8$$



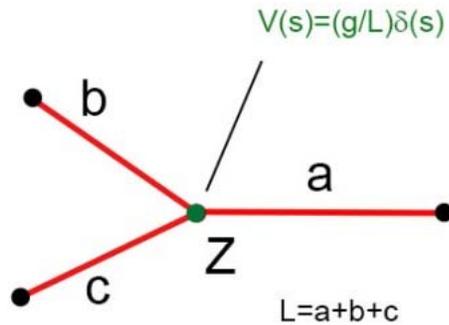
# Optimization of eigenstate overlap for max $\beta$ and $\gamma$



# Consistency check for 1,2, & 3 $\delta$ atoms



## 3-star delta graph



$$\begin{aligned}\phi_a(s) &= \frac{Z \sin k(a - s_a)}{\sin ka} \\ \phi_b(s) &= \frac{Z \sin k(b - s_b)}{\sin kb} \\ \phi_c(s) &= \frac{Z \sin k(c - s_c)}{\sin kc}\end{aligned}$$

### Secular function for star graph

$$F_{star}(a, b, c) = \frac{1}{4} [\cos k_n L_1 + \cos k_n L_2 + \cos k_n L_3 - 3 \cos k_n L]$$

$$L = a + b + c, \quad L_1 = |a + b - c|, \quad L_2 = |a - b + c|, \quad L_3 = |a - b - c|$$

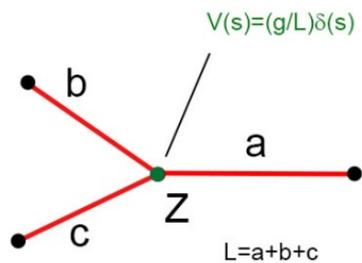
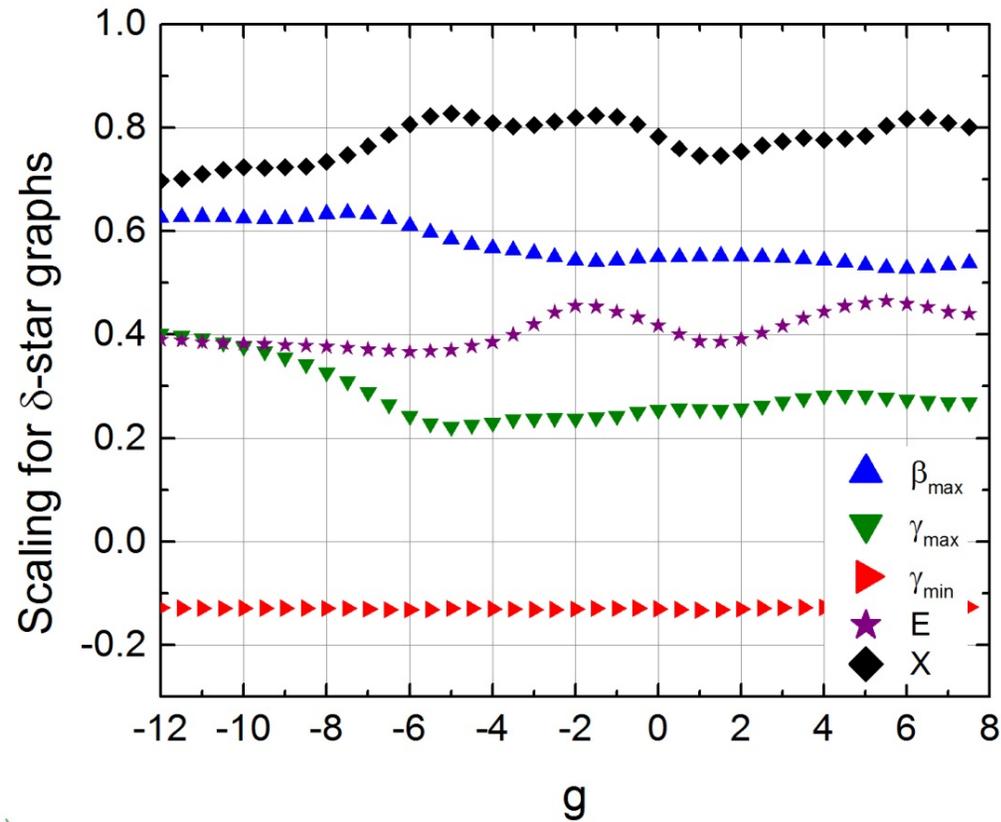
### Secular function for star delta graph

$$F_{star}(a, b, c) + (2g/kL) \sin ka \sin kb \sin kc = 0$$

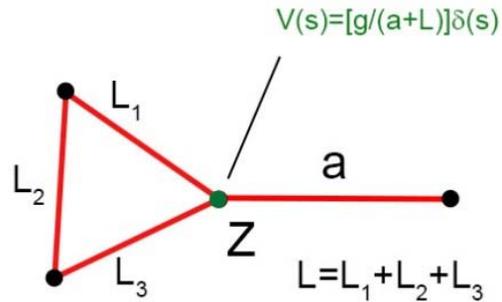
$$g_c = -\frac{(a + b + c)(ab + ac + bc)}{2abc}$$

$g < g_c$  yields one negative energy state

# Scaling of $\beta$ , $\gamma$ , $X$ , and $E$ for delta star graph



# Delta lollipop graph



$$\phi_a(s) = \frac{Z \sin k(a - s_a)}{\sin ka}$$

$$\phi_L(s) = \frac{Z \sin k(b - s_L) + Z \sin ks_L}{\sin kL}$$

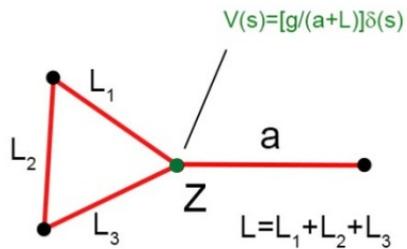
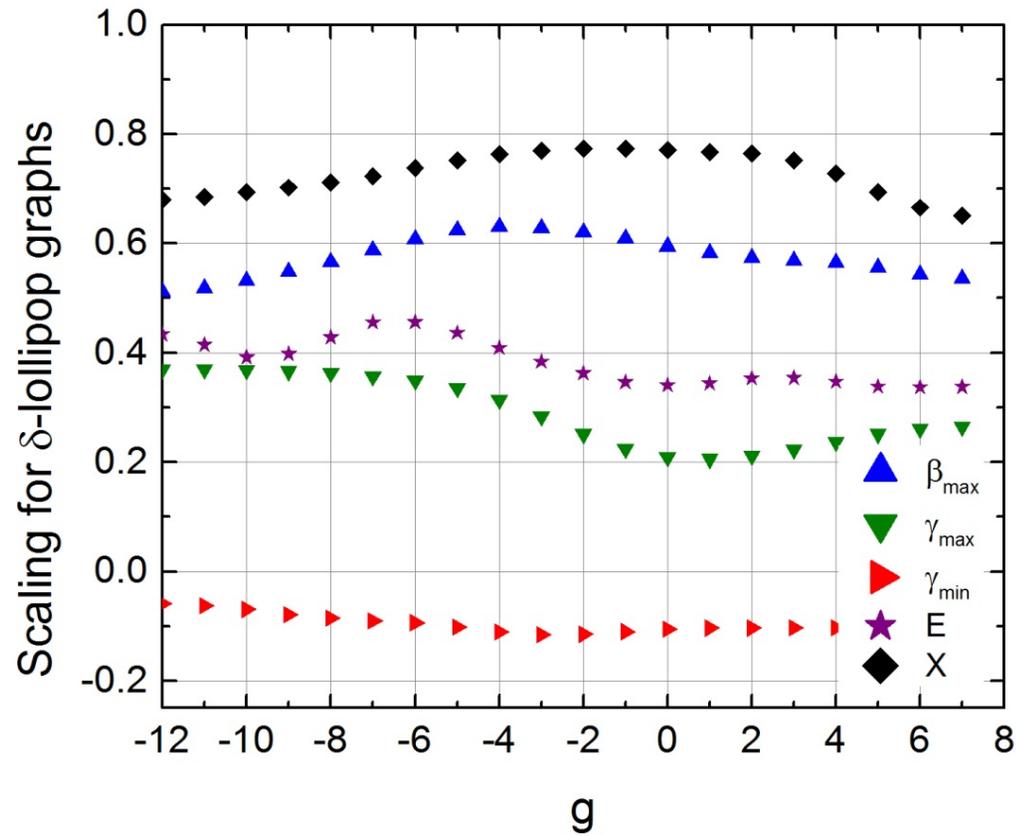
## Secular function for lollipop graph

$$F_{pop}(a, L) = \frac{1}{2} [3 \cos k(a + L/2) - \cos k(a - L/2)]$$

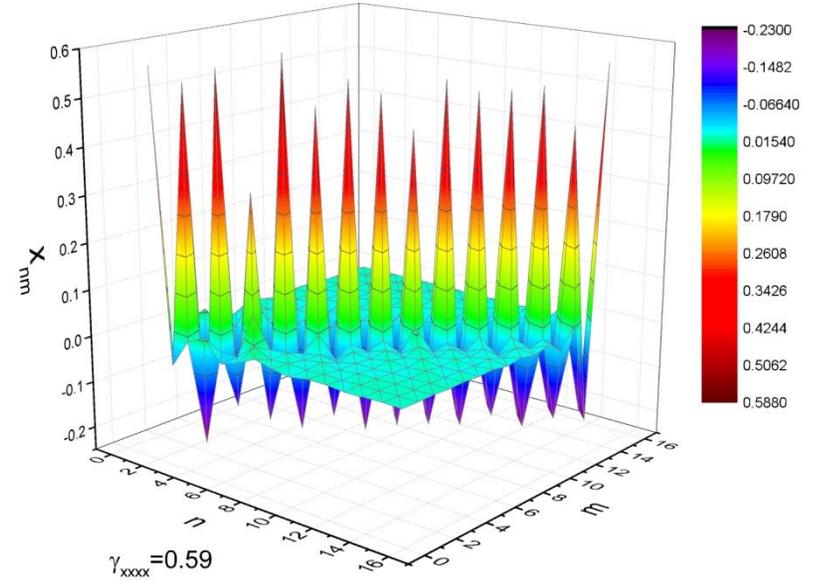
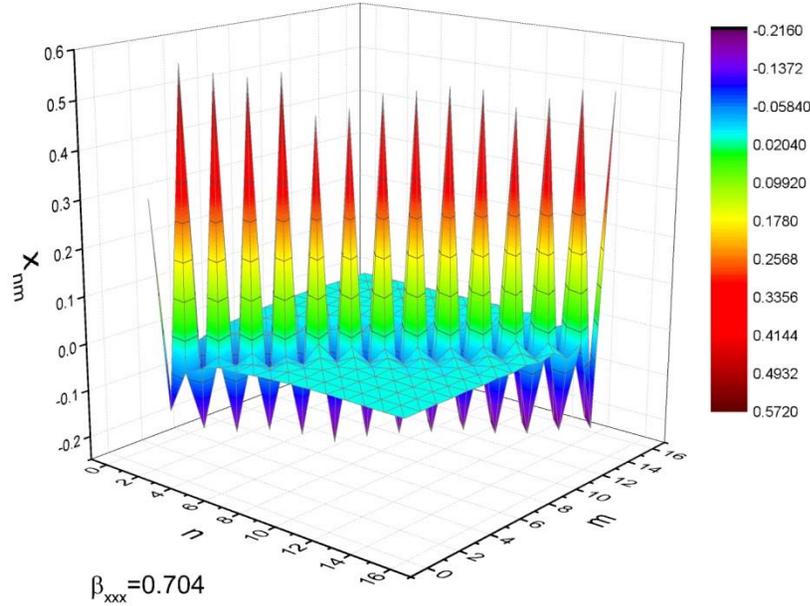
## Secular equation for lollipop delta graph

$$F_{pop}(a, L) + \frac{2g}{k(a + L)} \sin ka \cos kL/2 = 0$$

# Scaling of $\beta$ , $\gamma$ , $X$ , and $E$ for delta lollipop graph



# Transition moments for $3\delta$ graphs at max $\beta, \gamma$

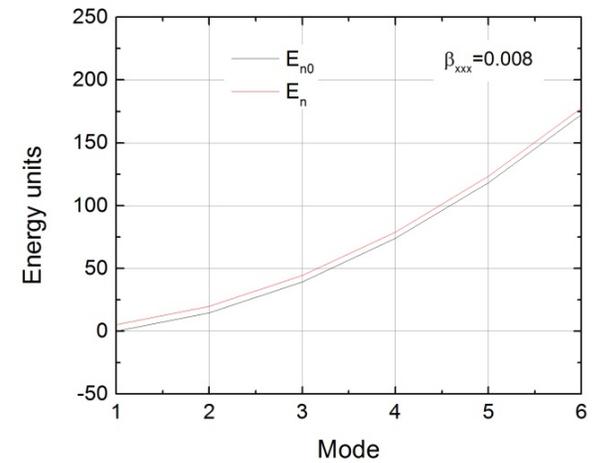
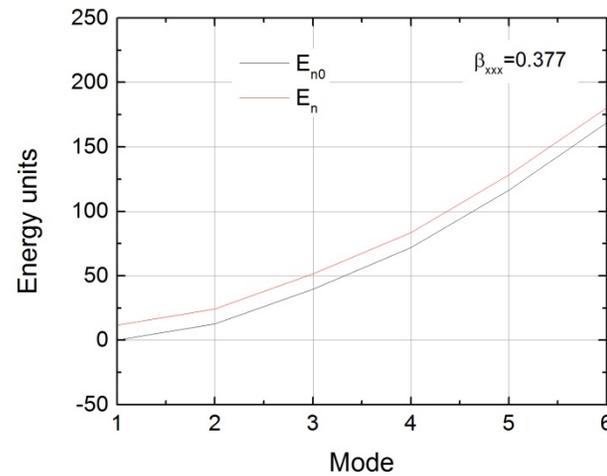
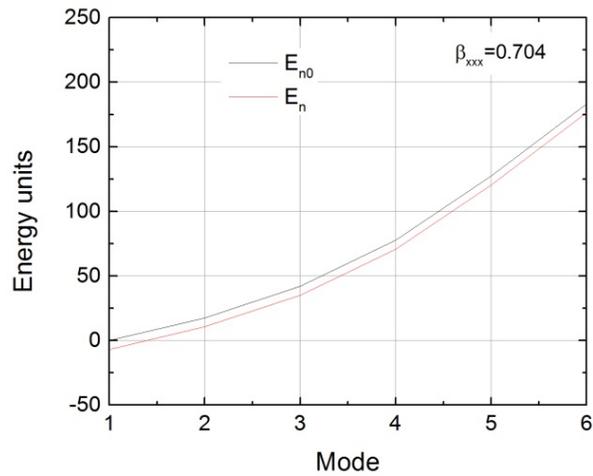
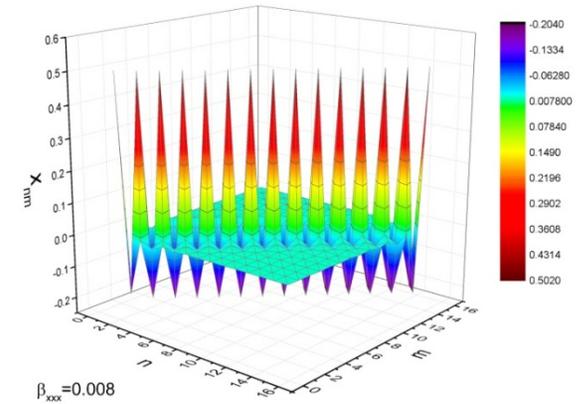
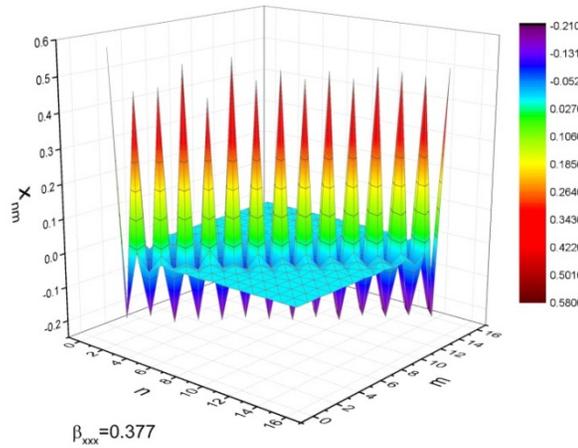
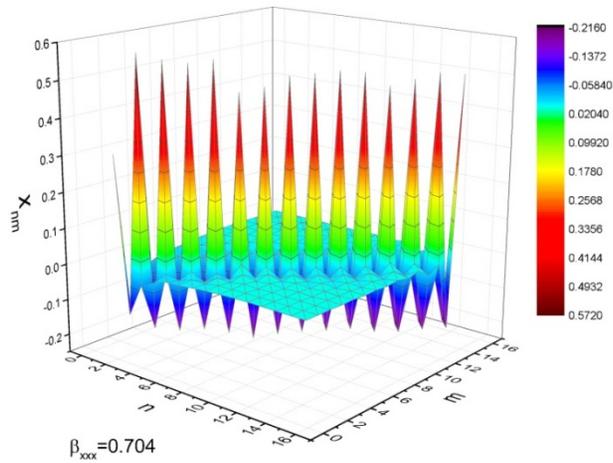


$$\beta_{xxx} = \sum'_{n,m} \beta_{xxx}(n, m) = 3^{3/4} E_{10}^{7/2} \sum'_{n,m} \frac{x_{0n} \bar{x}_{nm} x_{m0}}{E_{n0} E_{m0}}$$

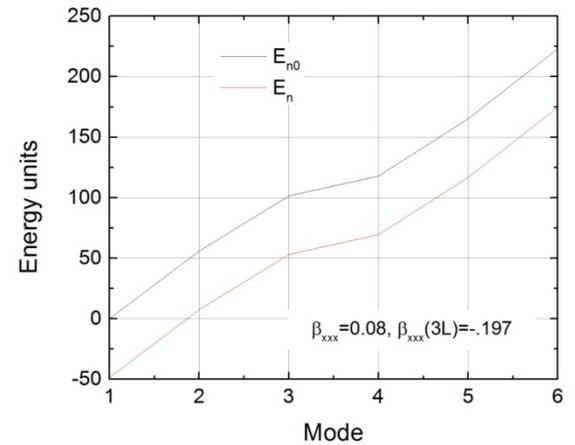
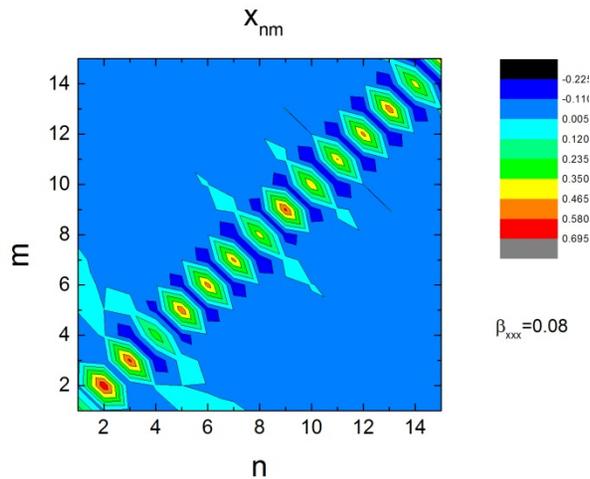
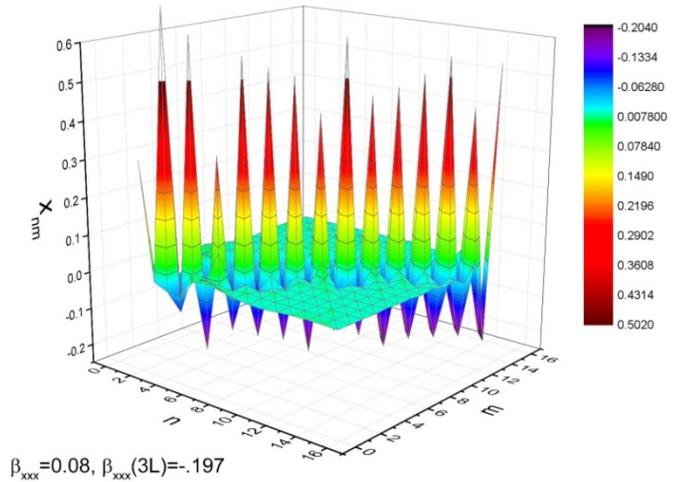
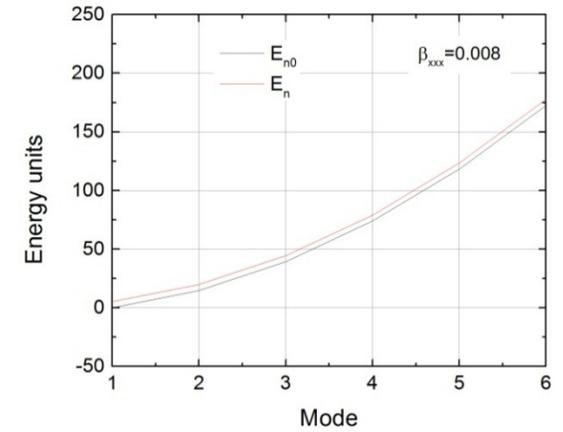
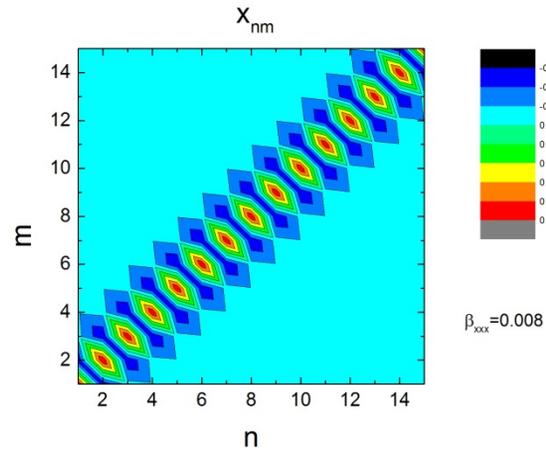
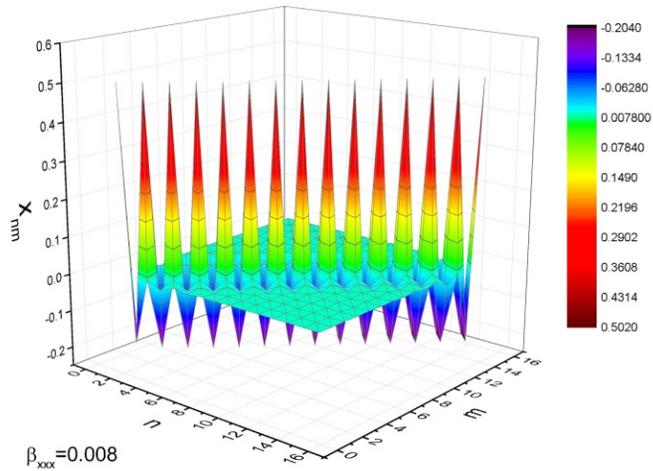
$$\beta_{xxx}^{3L} = 3^{3/4} E_{10}^{7/2} \left( \frac{|x_{01}|^2 \bar{x}_{11}}{E_{10}^2} + \frac{|x_{02}|^2 \bar{x}_{22}}{E_{20}^2} + \left[ \frac{x_{01} x_{20} x_{12}}{E_{10} E_{20}} + c.c. \right] \right)$$

$$\beta_{xxx}^{4L} = \beta_{xxx}^{3L} + 3^{3/4} E_{10}^{7/2} \left( \frac{|x_{03}|^2 \bar{x}_{33}}{E_{30}^2} + \left[ \frac{x_{03} x_{10} x_{31}}{E_{10} E_{30}} + c.c. \right] + \left[ \frac{x_{03} x_{20} x_{32}}{E_{20} E_{30}} + c.c. \right] \right)$$

# These are described by the 3L model

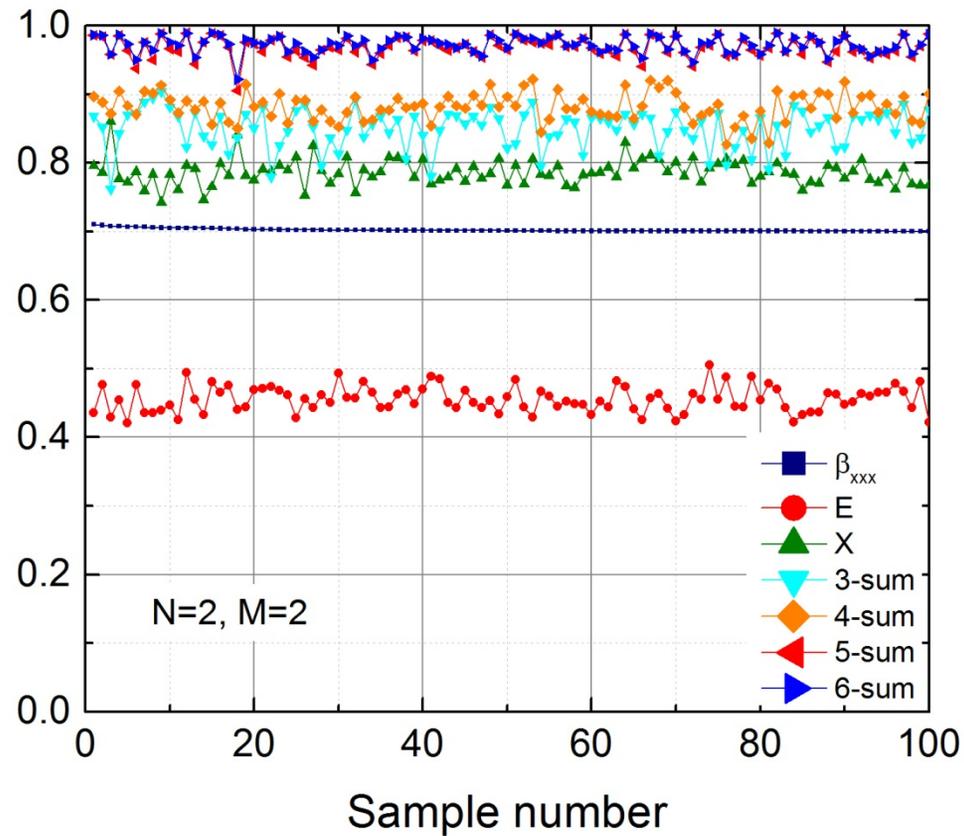


# 3L model for $\delta$ graphs with different $g$ but $\beta \sim 0$



$\beta_{xxx}(4L)=0.058$

# Partial sums when $\beta_{xxx}$ is near its maximum value



# Summary: For $\delta$ -dressed quantum graphs

Negative energy ground state significantly alters low-end spectrum,  
enabling larger contributions from low-lying states for small  $\beta, \gamma$

Scaling is nevertheless maintained near optimum values

Fine tuning of states and spectra possible in these graphs

Existence of  $\delta$ -graphs with  $\beta_{\max} \sim 0.71, \gamma_{\max} \sim 0.6, \gamma_{\min} \sim -0.15$

Saturation of values with number of  $\delta$  potentials  $\sim 3$ .